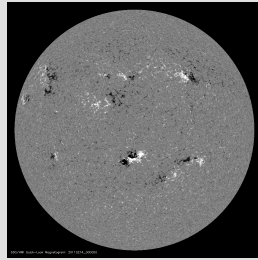
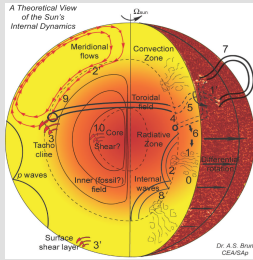


ASTR 7500: Solar & Stellar Magnetism

Hale CGEG Solar & Space Physics

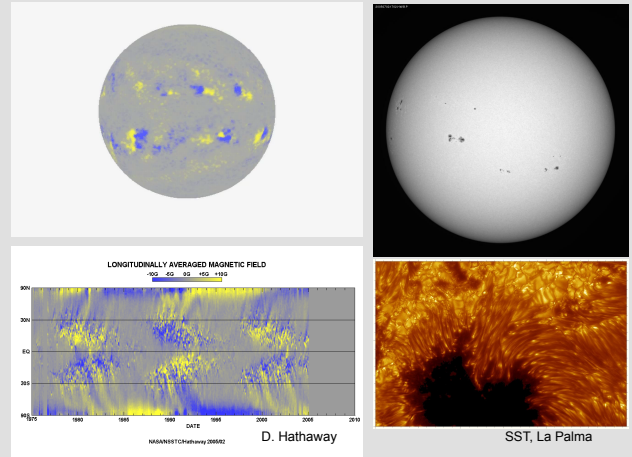


Matthias Rempel, Prof. Juri Toomre + HAO/NSO colleagues

Lecture 12 Thurs 28 Feb 2013

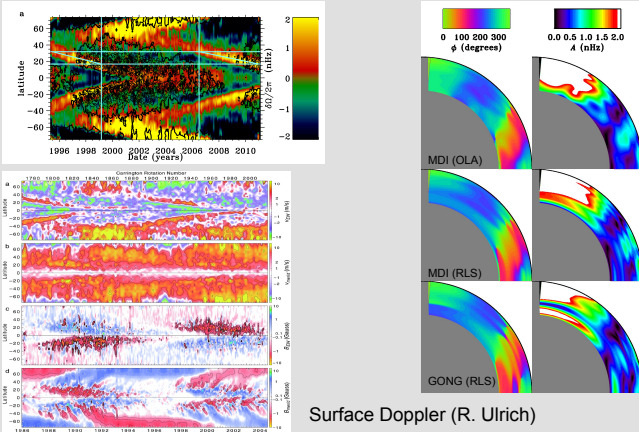
zeus.colorado.edu/astr7500-toomre

Solar magnetic field



Large scale flow variations

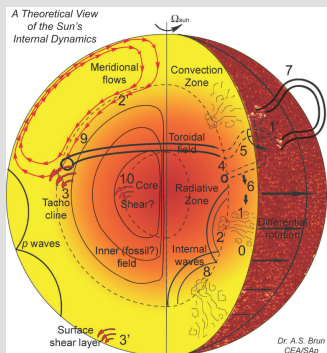
Global Helioseismology (R. Howe)



Solar dynamo models – what is the goal?

- What is a solar dynamo model supposed to do?
 - 1) Show a “solar-like” activity pattern in terms of:
 - Cyclic behavior with equator-ward propagation of activity
 - Surface flux evolution consistent with observations
 - Large scale flow variations consistent with observations
 - 2) Show a “solar-like” amplitude variation from cycle to cycle
 - 3) Allow prediction of future activity
- Most models struggle already with point 1)
 - Focus this lecture on 1)
 - 2) and 3) can provide additional constraints on dynamo models

The basic dynamo ingredients



- Large-scale flows
 - Differential rotation
 - Meridional flow
 - Turbulent and (cyclic) variation
- Turbulent induction
 - Transport
 - Advective
 - Diffusive
 - α -Effects
 - Key terms that enable dynamo action
- Flux emergence
 - Links dynamo to photospheric field observations
 - Might play role in dynamo process itself
 - Babcock-Leighton mechanism

Numerical modeling approaches

- Meanfield models
 - Solve equations for mean flows, mean magnetic field only
 - Inexpensive, but need good model for correlations of small scale quantities (e.g. turbulent angular momentum transport), see extensive work by Rüdiger & Kitchatinov
 - Can address the full problem, but not from first principles (models have many degrees of freedom and tunable parameters)
- 3D numerical simulations
 - Solve the full set of equations (including small and large scale flows, magnetic field) from first principles
 - Very expensive:
 - Low resolution runs for long periods >10 years
 - High resolution for short periods
 - Good understanding of differential rotation, ingredients of solar dynamo, no complete model yet
- Advances in computing infrastructure shift balance toward 3D simulations, but we need both!

Mean field models

- Mean field models consider only average quantities
 - Sunspots are a key feature of the solar cycle, but they are averaged away
- Mean field models make strong assumptions that are not well justified from first principles
- Too many degrees of freedom require “educated guesses”

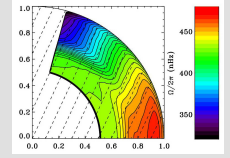
$$(\overline{v' \times B'})_i = a_{ik} \overline{B}_k + b_{ijk} \frac{\partial \overline{B}_j}{\partial x_k}$$

- Contains 36!!! (mostly unknown) functions of r and θ , in most models only 2 are considered and even that allows for a lot of freedom
- Computing mean field coefficients from 3D simulations (Schrinner et al. 2007, Ghizaru et al. 2011) shows that in general almost all of them are important!
- Mean field models allow us to study certain scenarios or they allow to analyze a complicated 3D simulation, but one has to be very lucky to find the “correct” model for the solar cycle without additional knowledge
- Non-linear feedback difficult to implement

Solar dynamo models

➤ Mean field models

- Convection zone dynamos
- Tachocline/interface dynamos
- Near surface shear layer dynamos
- Flux transport dynamos



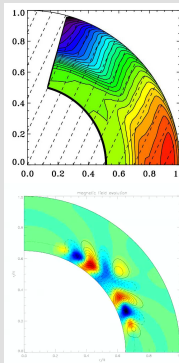
➤ Main uncertainties

- Location of dynamo
- Poloidal field regeneration (B_r , B_θ from B_ϕ : α -effect)
- Turbulent transport (magnetic pumping, turbulent diffusion vs. magnetic buoyancy)
- Role of meridional flow (propagation of activity belt)

Mean field dynamos

➤ Thin layer dynamos

- Overshoot/tachocline dynamos
 - Radial shear, $\alpha\Omega$ -type dynamos, latitudinal propagating dynamo wave
 - Negative α in northern hemisphere for equatorward propagation
- Surface shear layer?
- Main problem:
 - Typically very short latitudinal wave length (several overlapping cycles)



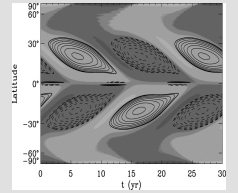
➤ Distributed dynamos

- Interface dynamos
 - Ω -effect in tachocline, α -effect in CZ, introduced to avoid problems with strong α -quenching
 - Solutions very sensitive to details

Mean field dynamos

➤ Distributed dynamos

- Flux transport dynamo
 - Advective transport of field by meridional flow
 - Propagation of AR belt advection effect
 - Cycle length linked to overturning time scale of meridional flow
- Central assumption:
 - Proper meridional flow profile (mostly single flow cell poleward at top, equatorward near bottom of CZ)
 - Weak turbulent transport processes
 - Babcock-Leighton α -effect
- Overall:
 - Most successful in reproducing solar like behavior



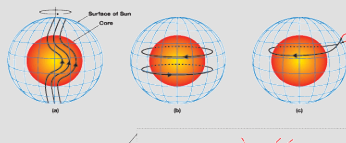
Dikpati et al. 2004

Schematic of a Babcock-Leighton flux transport model

(Durney, Choudhuri, Schüssler, Dikpati, Nandi, Charbonneau, Gilman, Rempel, Hotta)

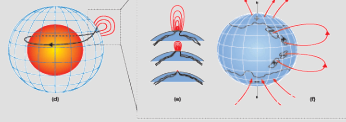
➤ Differential rotation

- Toroidal field production
- Stored at base of CZ
- Rising flux tubes

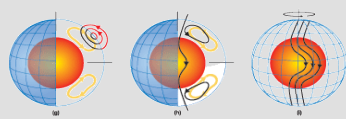


➤ Babcock-Leighton α effect

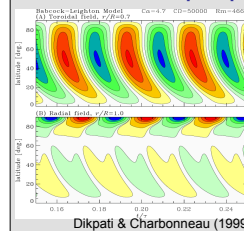
- Tilt angle of AR
- Leading spots have higher probability to reconnect across equator



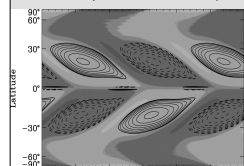
➤ Transport of magnetic field by meridional flow



Solution properties flux transport dynamos



Dikpati & Charbonneau (1999)



Dikpati et al. (2004)

➤ Good agreement with basic cycle properties

- Equatorward propagation
- Weak cycle overlap
- Correct phase relation between poloidal and toroidal field

➤ Less good agreement

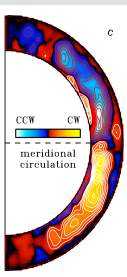
- Poleward extension of butterfly diagram?
- Polar surface field typically too strong
- Symmetry of solution (quadrupole preferred)

➤ More complicated ingredients can improve agreement

- Strong variation of magnetic diffusivity in CZ
- Additional α -effect at base of CZ

➤ Expense: Strong sensitivity to many not well known ingredients

Meridional flow structure, assumptions flux transport dynamo



3D simulation Miesch et al. (2008)

Observations

- Poleward near surface (surface Doppler and local helioseismology agree well)
- Recent results indicate shallow return flow (Hathaway 2011)?

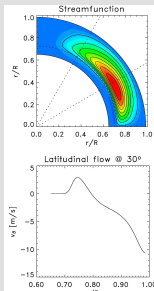
Theory

- Mean field models: single flow cell, related to inward transport of angular momentum
- 3D: low res runs multi cellular, recent high res single cell, results not yet converged

Advection dominated regime difficult to realize:

$$\eta_{turb} \propto H_p V_{rms}$$

$$V_{merid} \propto V_{rms}^2 / V_{rot}$$



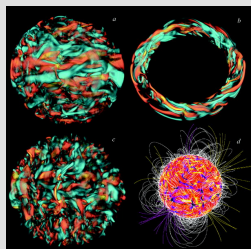
Mean field model Rempel (2005)

3D simulations

- Solve the full set of equations (including small and large scale flows, magnetic field) from first principles
 - No shortcuts, have to solve for the full problem including differential rotation and meridional flow
 - Non-linear effects automatically included
- Intrinsic limitations
 - Boundary conditions (radial direction)
 - Tachocline at base of CZ
 - Top boundary typically 20 Mm beneath photosphere
 - Cannot capture solar Re and Rm, how to treat small scales
 - DNS: resolve dissipation range with artificially increased diffusivities
 - (I)LES: do only the minimum required to maintain numerical stability
- Very expensive
 - Low resolution runs for long periods >10 years
 - High resolution for short periods
- Good understanding of differential rotation, ingredients of solar dynamo, no complete model yet

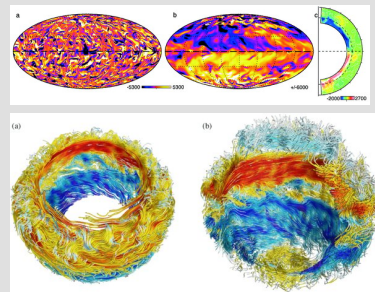
3D dynamo simulations

- 1981 Gilman & Miller
 - First 3D convective dynamos in a spherical shell (Boussinesq)
- 1983 Gilman
 - Dynamo simulations with reduced diffusivities
 - large scale field and periodic field reversal
 - poleward propagation
- 1985+ Glatzmaier ...
 - Mostly 3D geodynamo models
- 2004 Brun, Miesch, Toomre
 - Turbulent dynamo (anelastic)
 - 800 G peak toroidal field
 - Mean field 2% of energy
 - No cyclic behavior



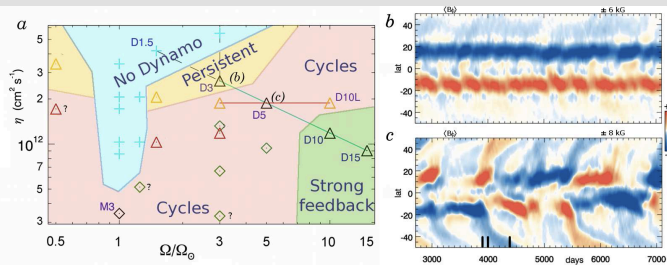
3D dynamo simulations

- 2006 Browning et al.
 - Addition of tachocline
 - Organized ~5 kG field in stably stratified region
- 2008+ Brown et al.
 - Faster rotating stars
 - Strong field (~10 kG) maintained within CZ
 - Cyclic behavior for certain parameter choices (faster rotation)



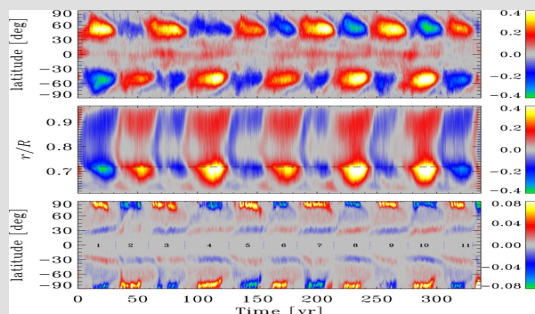
Cyclic dynamo regimes

- 2011 Brown et al.
 - Cyclic behavior typically found for sufficiently high Rm
 - Small diffusivity
 - Fast rotation

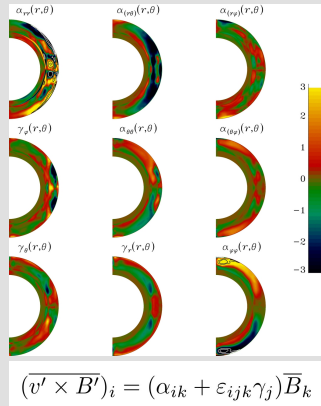


3D dynamo simulations

- 2010 Ghizaru et al., 2011 Racine et al.
 - Cyclic $\alpha\Omega$ -type dynamo 60 yr period
 - Magnetic field generated near base of CZ



Characterization of large-scale dynamo (Ghizaru et al. 2011)

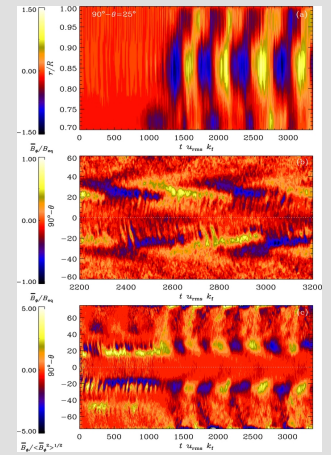


- All components have comparable amplitude
- Strongest effect present in diagonal elements, but substantial deviation from isotropy
 - $\alpha\theta\theta$ and $\alpha\phi\phi$ show pattern that reflects expectation from helicity profiles, but not arr
- Turbulent pumping
 - Down- and equatorward in bulk on convection zone
 - Poleward near surface in high latitudes
 - Mimics /competes with meridional flow

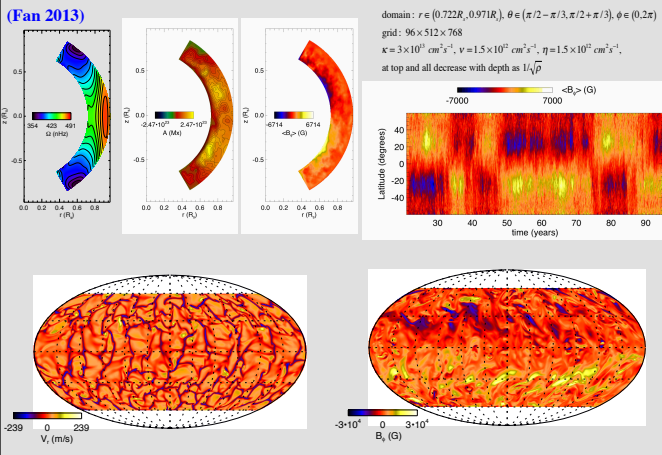
3D dynamo simulations

➤ Kapyla et al. (2012)

- 33 year period
- Field generated in bulk of CZ
- Equatorward propagation below 40 deg latitude
- Cycle length non-linear effect
 - Much shorter cycles during kinematic growth phase
 - “Phase transition” due to non-linear feedback



Simulations of convective dynamo in the solar convective envelope with the FSAM code



3D dynamo simulations

➤ Recent developments:

- Several independent groups find cyclic dynamos with periods in the 10-60 year range
- Some models with equatorward propagation of activity
- No simple explanation for cycle length and magnetic field patterns
 - Cycle length non-linear effect (longer cycles in saturated phase)
 - Not obvious if different models get similar solutions for the same reason

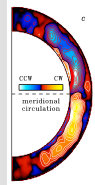
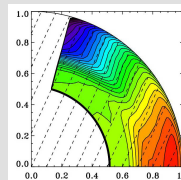
➤ Contrast to meanfield models:

- In general no single dominant turbulent induction term (like a scalar α -effect) that could capture the behavior
- Non-linear feedback more than just saturation effect (i.e. long cycle length only found in non-linear regime)

What are the main uncertainties?

➤ Large scale flows:

- Differential rotation well known
 - Role of latitudinal vs. radial shear not clear
 - Ω – effect: $B_p \cdot \nabla \Omega$
 - $\frac{\partial \Omega}{\partial r} > \frac{1}{r} \frac{\partial \Omega}{\partial \theta}$, but typically: $B_r < B_\theta$
- Role of tachocline (essential or does it just shape activity)
 - Fully convective stars show strong activity!
- Variation of Ω (torsional oscillations) very small
 - Weak magnetic feedback or DR strongly driven?
 - What does this tell us about saturation?
- Meridional flow
 - Poleward at surface
 - Flow structure in CZ?
 - Shallow return flow (Hathaway 2011)?



Miesch et al. 2008

What are the main uncertainties?

➤ Turbulent induction/transport

- In most 3D simulations turbulence is more complicated than a combination of diffusion, advection and α -effects
- Flux transport dynamos assume weak (< 10% of MLT estimates) turbulent transport processes - is that reasonable?
 - η has to be small, but not ν and κ (need to transport energy and maintain DR)?
 - no clear indication from numerical experiments for asymmetric magnetic quenching of η , ν and κ
- More general problem
 - Diffusivities of the order $\eta_{turb} \propto H_p V_{rms}$ give too short cycles
 - Are longer cycles an intrinsically non-linear effect?
- How is the poloidal magnetic field maintained?
 - kinematic (turbulent) α -effect?
 - magnetic saturation, role of magnetic helicity?
 - driven by magnetic instabilities?

What are the main uncertainties?

➤ Flux emergence process

- By-product of dynamo or essential part of dynamo process?
 - 10^{24} Mx is a lot of flux: $10 \text{ kG} \times 100 \text{ Mm}^2$
- Poloidal field in photosphere consequence of AR tilt angle
 - Babcock-Leighton α -effect
 - Is that enough to drive the dynamo?
 - Polar flux $\sim 0.1\%$ of toroidal flux
 - How to get back from 0.1% to 100%
 - DR can do $\sim 100!$
 - Babcock-Leighton flux transport dynamos have typically too strong polar field!

➤ What determines field amplitude

- Feedback on DR, meridional flow?
- Quenching of turbulent induction (magnetic helicity) ?

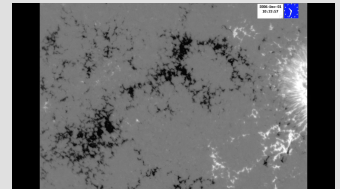
Flux emergence and sunspot formation

➤ General accepted view

- Magnetic flux rising toward surface from deep convection zone
- Observations show at first strong horizontal expansion of emerging flux

➤ Key question

- Transport of flux through convection zone and re-amplification in photosphere:
 - Density contrast of 10^6
 - $B \sim \rho^\epsilon$ $\epsilon = 1/2 \dots 2/3$
 - $100 \text{ kG} \rightarrow 100 \text{ G}$
 - $100 \text{ G} \rightarrow 3 \text{ kG} ???$
 - Vigorous convection



Flux emergence event observed with Hinode SOT

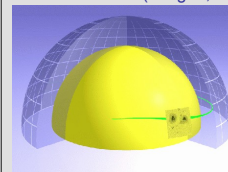
Modeling of flux emergence

➤ Lower convection zone (up to $\sim 20 \text{ Mm}$ depth beneath photosphere)

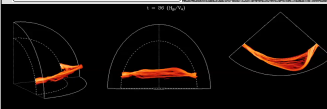
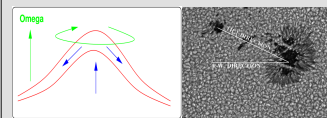
- Strongly subsonic velocities
- Ideal gas equation of state sufficient
- Size of flux tubes smaller than H_p and typical scale of convection
 - Flux tubes travel several times their diameter
 - Interaction with ambient flows (including flows created by rising flux) key to dynamics
 - Density contrast of 100 (out of 10^6)
- Modeling approaches
 - Thin flux tube approximation
 - 3D anelastic MHD models
 - Both with and without background convection

Flux emergence in lower convection zone

(Caligari, Fan, Fisher, Moreno-Insertis, Schüssler ...)



Thin flux tube simulation: Caligari et al. (1995)



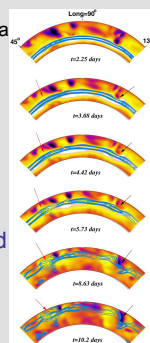
3D simulation: Fan (2008)

- Consistent results from thin tube and 3D simulations
- Coriolis force causes tilt of the top part of tube
- Explains asymmetry between leading/following spot
- Works best with $\sim 100 \text{ kG}$ flux tubes
 - Consistent with stability considerations in overshoot region
 - Too strong for dynamo models
- Twist required for 2D/3D simulations
 - Prevents fragmentation
 - Induces additional tilt (opposing that from Coriolis force)
 - Trade off between stability and tilt

Interaction with convection

➤ 3D anelastic MHD (Jouve & Brun 2009)

- Self consistent interaction with convection, differential rotation and meridional flow (Global convection zone simulation)
- Convective motions additional source of tilt, substantially shape tube during rise
- Challenge: Focus on global picture limits resolution on the scale of flux tube, requires tubes with $\gg 10^{22}$ Mx flux

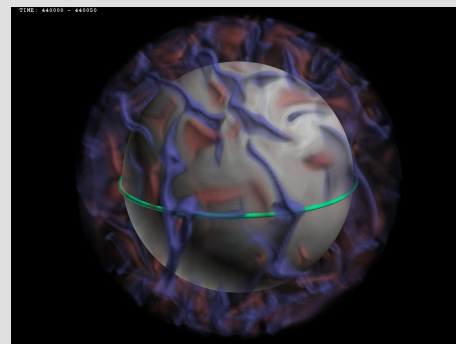


Jouve & Brun (2009)

➤ Thin flux-tubes rising in convective background

- Take velocity from global CZ simulation
- Treat flux tube as thin tube
- Weber, Fan & Miesch 2011

Interaction with convection



Weber, Fan & Miesch (2011)

- Thin flux tube rising in convective envelope (taken from global 3D simulation)
- Flux tube evolution mostly dominated by convective time scales
- Less dependence on initial field strength
 - Best results for $40\text{-}50 \text{ kG}$ ($100\text{-}150 \text{ kG}$ without convection)

Flux emergence in upper most 20 Mm

➤ Upper convection zone

- Subsonic/supersonic transition velocities
- Partial ionization, 3D radiative transfer important
- Size of flux tubes larger than H_p and typical scale of convection
 - Flux 'tubes' travel about their diameter
 - Density contrast of 10^4 (out of 10^6)
 - Dynamics dominated by strong expansion
 - Most weakening of field strength near surface
- Modeling approaches
 - Fully compressible MHD (with RT and realistic EOS)

➤ Currently treated independent from deep convection zone (computational constraints)

Flux emergence, sunspot formation

Magnetogram
($\tau=1$)

Domain size:
150x75x16 Mm

$t = 0.0$ h
30 Mm

$|B|$
(vertical slice)

Rempel & Cheung 2013

Sunspot fine structure

highest resolution, short time scales, shallow domains

Resolution:
16x16x12 km
(3072x3072x512)

Computing resource:
NSF-Teragrid
Cray-XT5 (Kraken, NICS)
3072-12288 cores
1 solar hour
~ 3 days (on 12288 CPUs)
~ 800,000 CPU hours
~ 10 TB data (3 TB/day !)

Data handling major challenge!
Achieved data rates larger than
SDO mission!

